

# Two-fluid magnetic reconnection with arbitrary guide-field

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The following new set of reduced equations governing 2-D, two-fluid, magnetic reconnection with arbitrary guide-field is derived:

$$\begin{aligned}\frac{\partial\psi_e}{\partial t} &= [\phi, \psi] + d_\beta [\psi, Z] + \eta \nabla^2 \psi, \\ \frac{\partial Z_e}{\partial t} &= [\phi, Z] + d_\beta [\nabla^2 \psi, \psi] + c_\beta [V_z, \psi] + c_\beta^2 \eta \nabla^2 Z, \\ \frac{\partial U}{\partial t} &= [\phi, U] + [\nabla^2 \psi, \psi] + \mu \nabla^2 U, \\ \frac{\partial V_z}{\partial t} &= [\phi, V_z] + c_\beta [Z, \psi] + \mu \nabla^2 V_z,\end{aligned}$$

where  $\psi$  is the magnetic flux-function,  $\phi$  the ion stream-function,  $Z$  the normalized perturbed magnetic field in the  $z$ -direction,  $U = \nabla^2 \phi$ ,  $\psi_e = \psi - d_e^2 \nabla^2 \psi$ , and  $Z_e = Z - c_\beta^2 d_e^2 \nabla^2 Z$ . Also,  $[A, B] \equiv \nabla A \wedge \nabla B \cdot \hat{\mathbf{z}}$ . Here,  $c_\beta = \sqrt{\beta/(1+\beta)}$ ,  $d_\beta = c_\beta d_i$ , where  $\beta$  is the plasma beta calculated with the guide-field,  $d_i$  the collisionless ion skin-depth, and  $d_e$  the collisionless electron skin-depth. It is assumed that  $\partial/\partial z \equiv 0$ . The reduction process involves the removal of the compressible Alfvén wave, which plays no role in magnetic reconnection.

The above equations, which are easy to solve numerically, represent a significant advance in magnetic reconnection theory, since all existing sets of reduced equations used to study reconnection are only valid in either the large or the small guide-field limits (*i.e.*, either  $\beta \gg 1$  or  $\beta \ll 1$ ). The new equations are used to investigate the scaling of the rate of forced magnetic reconnection in the so-called Taylor problem, in which a tearing stable plasma is subject to a suddenly applied resonant magnetic perturbation. In the non-linear regime with arbitrary guide-field (and  $d_e = 0$ ), the peak reconnection rate is independent of the resistivity, and scales like  $d\Psi/dt \sim [\beta/(1+\beta)]^{3/4} d_i^{3/2} \Xi_0^2$ , where  $\Xi_0$  the amplitude of the external perturbation. This result cannot be accounted for in terms of the ‘‘universal’’ reconnection rate recently proposed by the Maryland group.